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An Optimal Design Problem in Wave Propagation¹

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Abstract. In this paper we consider an optimal design problem in wave propagation proposed in (Ref. 1) in the one-dimensional situation: Given two materials at our disposal with different Young's modulus and different density, the problem consists of finding the best distributions of the two initial materials in a rod in order to minimize the vibration energy in the structure under periodic loading of driving frequency Ω . We comment on relaxation and optimality conditions, and perform numerical simulations of the optimal configurations. We also prove the existence of classical solutions in certain cases.

Key Words. Optimal design, wave propagation, classical solutions.

1 Introduction

In (Ref. 1) a structural model was proposed for the design of finite structures made up of two given materials in the context of wave propagation. The model is an optimal design problem in which the distribution of two given materials is optimized so as to minimize a cost functional related to the vibration or propagation of waves along the medium. In practice, this model may be useful for the systematic design of wave filters, damping of waves, or wave guides. The authors develop a numerical method for the optimization of those structures based on topology optimization (see (Ref. 2)), but there is no mathematical analysis of the model. In the numerical examples, they also observe the surprising fact that no microstructure appears between the two materials when one tries to design a filter or to minimize the vibration energy. Our aim here is to analyze mathematically the model proposed in (Ref. 1) in the one-dimensional situation for longitudinal propagation.

We consider the one-dimensional wave equation

$$\rho w_{tt} = (E w_x)_x,$$

where t is the time variable, x the space variable, w is the displacement, and E and ρ are the Young's modulus and material density respectively. Notice we do not consider any damping term in our equation. We are interested in time-harmonic solutions under periodic loading of the previous equation, i.e. in solutions of the form

$$w(x, t) = u(x) \exp(i\Omega t),$$

where Ω is the driving frequency and u the amplitude function. Then the amplitude, u , verifies the equation

$$(E(x)u')' + \Omega^2 \rho(x)u = 0.$$

Now, let us assume that we have two materials at our disposal, and consequently the elastic coefficient, $E(x)$, and the density, $\rho(x)$, take two values on, E_1, E_2 ($0 < E_1, E_2$) and ρ_1, ρ_2 ($0 < \rho_1, \rho_2$) respectively. We make the technical assumption

$$E_1 \leq E_2, \quad \rho_1 \leq \rho_2,$$

and that at least one of the inequalities is strict (otherwise we would just have one material). If we have a rod of length l , that we represent by the interval $[0, l]$, and we

distribute the two materials in the rod, the elastic coefficient function and the density function are respectively

$$E(x) = E_1\chi(x) + E_2(1 - \chi(x)), \quad \rho(x) = \rho_1\chi(x) + \rho_2(1 - \chi(x)),$$

where χ is the characteristic function of the (measurable) subset of the rod where we place the first material. Now it is natural to consider the following optimal design problem,

$$(P) \quad \text{Min}_{\mathcal{A}_{ad}} \quad J(\chi),$$

where the admissible set of designs is

$$\mathcal{A}_{ad} = \{\chi : \text{characteristic function of a measurable subset of } [0, l]\},$$

the cost functional is the vibration energy

$$J(\chi) = \int_0^l [\Omega^2(\rho_1\chi(x) + \rho_2(1 - \chi(x)))|u(x)|^2 + (E_1\chi(x) + E_2(1 - \chi(x)))|u'(x)|^2] dx,$$

and u is computed from χ through the state equation

$$[(E_1\chi(x) + E_2(1 - \chi(x)))u']' + \Omega^2(\rho_1\chi(x) + \rho_2(1 - \chi(x)))u = 0,$$

$$(E_1\chi(0) + E_2(1 - \chi(0)))u'(0) = \gamma, \quad (E_1\chi(l) + E_2(1 - \chi(l)))u'(l) = 0,$$

γ is a fixed given nonzero number. Here we consider the case of boundary conditions corresponding to one external load on the left extreme of the rod, and homogeneous Neumann boundary conditions on the other extreme. Roughly speaking, the optimal design problem consists of determining the best distribution (i.e., the best χ) of the two given materials in the rod in order to minimize the vibration energy in the structure, so that we minimize the vibration amplitude along the rod. Notice that our optimal design problem indeed is an optimal design problem for an ODE, as we have simplified the wave equation by considering time harmonic solutions.

We do not consider any volume constraint in this problem, and this is in accordance with the physical nature of the problem since waves propagate better through homogeneous materials than through mixtures of two materials. However the problems makes perfect sense if we impose a restriction on the amount of the materials that we may use, as for instance the constraint

$$\int_0^l \chi(x) dx \leq \delta, \text{ or } \int_0^l \chi(x) dx = \delta,$$

with $\delta \in (0, l)$ a fixed constant.

The solutions of the state equations are understood as weak solutions. That equation is obviously nonelliptic; however by the Fredholm's alternative, we can claim the existence of a unique solution for any boundary data if and only if Ω^2 is not an eigenvalue for the problem

$$-[(E_1\chi(x) + E_2(1 - \chi(x)))u']' = \lambda(\rho_1\chi(x) + \rho_2(1 - \chi(x)))u, \quad u \in H^1(0, l), \quad (1a)$$

$$(E_1\chi(0) + E_2(1 - \chi(0)))u'(0) = 0, \quad (1b)$$

$$(E_1\chi(l) + E_2(1 - \chi(l)))u'(l) = 0. \quad (1c)$$

Recall that, as the operator involved is compact and self-adjoint, we know that there exists an increasing (maybe not strictly) sequence of eigenvalues. For the sake of simplicity, and in order to formulate the optimal design problem, we will assume that Ω^2 is far away of any eigenvalue of the problem (1a-1c) for any characteristic function χ . This hypothesis makes physical sense (see (Ref. 1)).

It is well-known that in principle we cannot hope the previous optimal design problem

to admit optimal solutions, and in general we need to study relaxation for it (see (Ref. 8) and the references therein).

The plan of the paper is the following. Section 2 is devoted to state a relaxation result for our optimal design problem and to give necessary conditions of optimality for the relaxed formulation of the problem. In Section 3 we obtain qualitative properties of the optimal solutions of the relaxed problem, being able to prove, in certain cases, that there are optimal solutions in the form of characteristic functions, what implies the existence of minimizers for the original optimal design problem (P). Finally, in Section 4 we perform numerical simulations by using a gradient method based on the sensitivities computed in Section 2 and analyze numerically examples of practical interest. These examples are striking, as we can see how the vibration amplitude can be almost completely damped in certain situations by the design of the structure of the rod, and this shows how rich the behavior of the system is under the design of its internal structure, having into account that there is no damping in our model.

We would like to finish this introduction surveying about some problems related to

the optimal design problem subject of this paper, either in a physical or mathematical sense, or both of them. The problem most related to our optimal design problem is the maximization of band-gaps. Band gap materials prevent waves of frequencies belonging to a certain interval (the band gap interval) to propagate through the material. In this context, the optimal design problem of maximizing the band gap of a medium by optimizing the distribution of two given phases arises naturally. The pioneering papers (Refs. 4-5) (see also (Ref. 1), where the elastic case is considered) analyze this problem from a mathematical and numerical perspective. This is an optimal design problem in which we have to design a unit cell of a infinite periodic medium finding the best distribution of two (given) materials in order to maximize the band-gap interval. Recall that if a wave has frequency of propagation inside the band-gap, then it cannot propagate through the material, so that we are designing the material in order to avoid the propagation of waves. If you have a certain annoying frequency Ω that you would like to filter then you can create a medium (or material), with a periodic structure, such that Ω belongs to a band-gap interval of such a medium. Rather than this, the aim of our problem is to

design a finite structure (instead of a periodic infinite medium) in order to minimize the vibration energy of waves of driving frequency Ω . There is no periodicity constraint on the designs of our problem.

There are very few examples in the literature of optimal control problems in the coefficients for wave equations. In (Ref. 6) the authors consider an optimal control problem for the scalar linear wave equation, where the control variable is the sound speed which is assumed to be a function of time exclusively. A more general problem is studied in (Ref. 7), in which the design variables are the coefficients of the wave equation, and those are considered as functions of both space and time, see also the survey (Ref. 8) on the theory of dynamic materials.

2 Relaxation

Relaxation for optimal control or optimal design problems for ODE's is a classical subject. In this sense our optimal design problem fits into this framework and relaxation

for it is straightforward, and so to prove the relaxation result stated below is standard (basic references on the subject are (Ref. 8, 9, 10)). Another way to prove the result is by using H -convergence (Ref. 12) and G -closures techniques (Ref. 11), those were developed for the PDE's context and are applicable on our situation, although unnecessary.

Theorem 2.1 The optimal design problem

$$(\tilde{P}) \quad \text{Min}_{\mathcal{A}_{ad}^*} \quad \tilde{J}(E(\theta), \rho(\theta)),$$

where the admissibility set is

$$\mathcal{A}_{ad}^* = \{\theta \in L^\infty(0, l) : 0 \leq \theta \leq 1\},$$

$E(\theta)$ and $\rho(\theta)$ are given by

$$E = E(\theta) = E_1 E_2 / (E_1 + \theta(E_2 - E_1)), \quad \rho = \rho(\theta) = \rho_1 \theta + \rho_2 (1 - \theta),$$

the cost functional is

$$\tilde{J}(E, \rho) = \int_0^l [\Omega^2 \rho(x) |u(x)|^2 + E(x) |u'(x)|^2] dx,$$

and u is computed from the pair (E, ρ) through the state equation

$$[Eu']' + \Omega^2 \rho u = 0, \quad (2a)$$

$$E(0)u'(0) = \gamma, \quad E(l)u'(l) = 0 \quad (2b)$$

is a relaxation of the optimal design problem (P).

Notice that, in the case we consider a volume constraint, in our optimal design problem, Theorem 2.1 is still true. We will need for the numerical simulations the gradient of the cost functional \tilde{J} in the relaxed problem (\tilde{P}) . Its computation is direct by using the classical adjoint method.

Theorem 2.2 The functional \tilde{J} in problem (\tilde{P}) is Gâteaux differentiable on the set \mathcal{A}_{ad}^* , and the Gâteaux derivative at $\bar{\theta}$ in the direction of the admissible variation θ (admissible variations means that $\theta \in L^\infty(0, l)$ such that there exists $h_0 > 0$ such that for every

$h \leq h_0$, $\bar{\theta} + h\theta \in \mathcal{A}_{ad}^*$) is given by

$$\tilde{J}'(\bar{\theta}; \theta) = \int_0^l \{ [E'(\bar{\theta}(x))(\bar{u}'(x)\bar{p}'(x) + |\bar{u}'(x)|^2)$$

$$+ \Omega^2(\rho_2 - \rho_1)(\bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2)]\theta(x)$$

$$- (E'(\bar{\theta}(x))\bar{u}(x)\bar{p}(x)\theta(x))'\} dx,$$

where \bar{u} is the state associated to $\bar{\theta}$ (i.e. the solution of (2a-2b) for the pair $(E(\bar{\theta})\rho(\bar{\theta}))$),

and \bar{p} is the adjoint state, being in this case the unique solution of

$$[E(\bar{\theta})\bar{p}]' + \Omega^2\rho(\bar{\theta})\bar{p} = 2(\Omega^2\rho(\bar{\theta})\bar{u} - [E(\bar{\theta})\bar{u}]'), \quad (3a)$$

$$E(\bar{\theta}(0))\bar{p}'(0) = -2\gamma, \quad E(\bar{\theta}(l))\bar{p}'(l) = 0. \quad (3b)$$

Using the Taylor expansion of the cost functional we obtain, as a direct consequence of the previous result, the following maximum principle.

Corollary 2.1 If $\bar{\theta}$ (with $\bar{\theta}$ an element of the interior of \mathcal{A}_{ad}^*) is an optimal solution for

problem (\tilde{P}) then

$$\tilde{J}'(\bar{\theta}; \theta - \bar{\theta}) \geq 0$$

for any θ such that $\theta - \bar{\theta}$ is an admissible variation.

3 Existence of Solutions for Problem (P)

In this section we explore the consequences of the optimality conditions given in the form of a maximum principle in Corollary 2.1 on the optimal solutions. Our objective is to check whether it is possible to prove existence results for the original optimal design problem (P). In the case that $E_1 = E_2$ (and we assume then that $E_1 = E_2 = 1$), the design variable does not act on the principle part of the state equation and the problem is mathematically much simpler, we are able to prove that indeed any optimal solution for problem (\tilde{P}) is a characteristic function, so that the original optimal design problem admits minimizers. We use the optimality conditions in this particular case to show the result, which is given in Theorem 3.1. The general situation is more difficult and does

not seem possible to use the same argument, however we give a condition implying the existence of at least one optimal solution in the form of a characteristic function. We divide this section into two subsections corresponding to the two cases.

3.1 Case $E_1 = E_2$

We assume that $E_1 = E_2 = 1$ and our problem now, that we denote as (\tilde{P}_1) , is the following

$$\text{Min}_{\mathcal{A}_{ad}^*} \tilde{J}(\rho(\theta)),$$

where the admissibility set is

$$\mathcal{A}_{ad}^* = \{\theta \in L^\infty(0, l) : 0 \leq \theta \leq 1\},$$

$\rho(\theta)$ is given by

$$\rho(\theta) = \rho_1\theta + \rho_2(1 - \theta),$$

the cost functional is

$$\tilde{J}(\theta) = \int_0^l [\Omega^2 \rho(\theta(x)) |u(x)|^2 + |u'(x)|^2] dx,$$

and u is computed from θ through the state equation

$$[u']' + \Omega^2 \rho(\theta(x))u = 0,$$

$$u'(0) = \gamma, \quad u'(l) = 0$$

Following Theorem 2.1, this problem is in the relaxed form and therefore it admits at least one optimal solution. The following result establishes that any optimal solution is actually a classical one (classical solution means a solution in the form of a characteristic function) .

Theorem 3.1 Any optimal solution, $\bar{\theta}$, for problem (\tilde{P}_1) takes on only the values $\{0, 1\}$, i.e. any optimal solution for (\tilde{P}_1) is a characteristic function.

Proof:

Using Theorem 2.2 we have that

$$\tilde{J}'(\bar{\theta}; \theta) = \int_0^l \Omega^2(\rho_2 - \rho_1)\theta(x) [\bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2] dx,$$

where \bar{u} is the state associated to $\bar{\theta}$, i.e. the solution of

$$[\bar{u}']' + \Omega^2 \rho(\bar{\theta}) \bar{u} = 0,$$

$$\bar{u}'(0) = \gamma, \quad \bar{u}'(l) = 0$$

and \bar{p} is the adjoint state, in this case \bar{p} is the unique solution of the equation

$$[\bar{p}']' + \Omega^2 \rho(\bar{\theta}) \bar{p} = 2 (\Omega^2 \rho(\bar{\theta}) \bar{u} - \bar{u}''),$$

$$\bar{p}'(0) = -2\gamma, \quad \bar{p}'(l) = 0.$$

By Corollary 2.1 we obtain the maximum principle

$$\int_0^l \Omega^2 (\rho_2 - \rho_1) (\theta(x) - \bar{\theta}(x)) [\bar{u}(x) \bar{p}(x) - |\bar{u}(x)|^2] dx \geq 0 \quad (4)$$

for any admissible variation θ . Now choosing θ as

$$\theta(x) = \begin{cases} 1 & \text{if } \bar{u}(x) \bar{p}(x) - |\bar{u}(x)|^2 < 0, \\ \bar{\theta}(x) & \text{if } \bar{u}(x) \bar{p}(x) - |\bar{u}(x)|^2 = 0, \\ 0 & \text{if } \bar{u}(x) \bar{p}(x) - |\bar{u}(x)|^2 > 0, \end{cases}$$

for a.e. $x \in [0, l]$, we have that the integrand in (4) is less or equal than zero, but the

integral is positive, and thus the only possibility is that the integrand vanishes at almost

every point, and therefore

$$\bar{\theta}(x) = 1 \quad \text{if } \bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2 < 0,$$

$$\bar{\theta}(x) \in (0, 1) \quad \text{if } \bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2 = 0,$$

$$\bar{\theta}(x) = 0 \quad \text{if } \bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2 > 0,$$

for a.e. $x \in [0, l]$. Now we will prove that

$$\bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2 \neq 0$$

a.e. $x \in [0, l]$.

Notice that $\bar{u} \in H^2(0, l)$ and then, by the Sobolev embedding, $u \in C^1(0, l)$. By the existence and uniqueness of solutions for initial value problems for ODE's (applied to u), and the identity

$$\bar{u}(x)\bar{p}(x) - |\bar{u}(x)|^2 = \bar{u}(x) (\bar{p}(x) - \bar{u}(x))$$

we just need to prove that $v(x) = \bar{p}(x) - \bar{u}(x)$ vanishes on a zero measure set of $[0, l]$, but

this is a direct consequence of the fact that v verifies the equation

$$v'' + \Omega^2 \rho(\bar{\theta})v = 2 (\Omega^2 \rho(\bar{\theta})\bar{u} - \bar{u}'') = 4\Omega^2 \rho(\bar{\theta})\bar{u}$$

in which the right hand side is different of zero a.e. $x \in [0, l]$. Therefore we conclude that

$\bar{\theta}$ only takes on the values 0 and 1. □

An obvious observation is that there is no design such that the corresponding amplitude is extinct in finite time.

3.2 General case

The existence of classical solutions for the general situation of (\tilde{P}) is not clear in general as far as the authors are able to claim, although something may be said about this. The rest of the section is devoted to discuss this issue. We do not state a formal result in this section, but show using new ideas that the problem admits optimal solutions in certain cases. We think this discussion is interesting.

In Theorem 3.1 we have shown that any optimal solution for problem (\tilde{P}_1) is classical, but for our proposes it would be also interesting just to know that there exists at least a classical solution. In order to explore this for problem (\tilde{P}) we proceed in the following way. We know, by Theorem 2.1, that there exists at least one optimal design, $\bar{\theta}$, for

problem (\tilde{P}) . Let us call

$$\bar{\delta} = \int_0^l \bar{\theta}(x) dx,$$

and as the volume fraction is not constrained in our optimal design problem, we may say

that $\bar{\delta}$ is an optimal volume fraction. Now we would like to see whether there exists a

characteristic function $\bar{\chi}$ such that

$$\bar{\delta} = \int_0^l \bar{\chi}(x) dx, \quad \tilde{J}(\bar{\theta}) = \tilde{J}(\bar{\chi}),$$

so that $\bar{\chi}$ is an optimal design for problem (\tilde{P}) , and therefore it admits at least one classical

optimal solution. To this end, it is useful to introduce the following standard change of

variables, for a given θ with $\int_0^l \theta(x) dx = \bar{\delta}$, set

$$y = \int_0^x (E_1 + \theta(E_2 - E_1)) / (E_1 E_2) dx,$$

then defining $U(y) = u(x)$, it is very easy to check that the state equation for the optimal

design problem (\tilde{P}) is reduced to

$$U''(y) + \Omega^2 \Psi(y) U(y) = 0, \quad y \in [0, L]$$

where $L = y(l)$ (notice that L only depends on the volume fraction of θ , i.e. $\bar{\delta}$, which is fixed), the coefficient of the equation

$$\Psi(y) = E_1 E_2 (\rho_2 + \theta(x)(\rho_1 - \rho_2)) / (E_1 + \theta(x)(E_2 - E_1))$$

and the cost functional after the change of variables is

$$I(\Psi(y)) = \tilde{J}(\rho(x)) = \int_0^L [\Omega^2 \Psi(y) |U(y)|^2 + |U'(y)|^2] dy.$$

The change of variables that we use was introduced in (Ref. 13) for a related problem of optimization for eigenfrequencies.

It is elementary to check that function f ,

$$f(t) = E_1 E_2 (\rho_2 + t(\rho_1 - \rho_2)) / (E_1 + t(E_2 - E_1)),$$

attains its extreme points at the extreme points of the interval $[0, 1]$, and therefore

$$E_1 \rho_1 \leq \Psi(y) \leq E_2 \rho_2,$$

a.e. $y \in [0, L]$. Actually, since function f is one-to-one, the transformation $f(\theta)$ is a one-to-one transformation of \mathcal{A}_{ad}^* onto

$$\{\Psi \in L^\infty(0, L) : E_1 \rho_1 \leq \Psi(y) \leq E_2 \rho_2\}.$$

The function Ψ verifies trivially the volume constraint

$$\int_0^L \Psi(y) dy = \rho_2 + \bar{\delta}(\rho_2 - \rho_1) = \delta$$

as the first moment of $\bar{\theta}$ is fixed and equal to $\bar{\delta}$.

Summarizing, what we have got with the change of variables is to prove that problem

(\tilde{P}) subject to the additional volume constraint

$$\bar{\delta} = \int_0^l \theta(x) dx,$$

is equivalent to the problem

$$\text{Min}_{\mathcal{A}_{ad}^*} I(\Psi(y)) = \int_0^L [\Omega^2 \Psi(y) |U(y)|^2 + |U'(y)|^2] dy$$

where

$$\mathcal{A}_{ad}^* = \left\{ \Psi \in L^\infty(0, L) : E_1 \rho_1 \leq \Psi(y) \leq E_2 \rho_2, \int_0^L \Psi(y) dy = \delta \right\}$$

and $U(y)$ is the unique solution of

$$U''(y) + \Omega^2 \Psi(y) U(y) = 0, \quad y \in [0, y(l)],$$

$$U'(0) = \gamma, \quad U'(L) = 0$$

If we call $\sigma_1 = E_1\rho_2$, $\sigma_2 = E_2\rho_2$, and $\sigma(t) = \sigma_1 + t(\sigma_2 - \sigma_1)$, this optimal design problem can be reformulated in the following way:

$$\text{Min}_{\mathcal{U}_{ad}^*} \quad I(\theta) = \int_0^L [\Omega^2 \sigma(\theta(y)) |U(y)|^2 + |U'(y)|^2] dy$$

where

$$\mathcal{U}_{ad}^* = \{\theta \in L^\infty(0, L) : 0 \leq \theta(y) \leq 1\}$$

and $U(y)$ is the unique solution of

$$U''(y) + \Omega^2 \sigma(\theta(y)) U(y) = 0, \quad y \in [0, L],$$

$$U'(0) = \gamma, \quad U'(L) = 0$$

and subject to the volume constraint

$$\delta = \int_0^L \theta(y) dy. \tag{5}$$

This is a problem exactly of the kind of problem (\tilde{P}_1) but subject to the additional volume constraint (5). As we mentioned at the end of Section 2, it is straightforward to conclude that the problem admits at least an optimal solution in the class \mathcal{U}_{ad}^* satisfying the constraint (5), as the constraint is weak- \star continuous in $L^\infty(0, L)$. Indeed we can go

further, and if we introduce a Lagrange's multiplier for the volume constraint (5), the results of Section 2 still hold and we can reproduce the computations of Theorem 3.1, to obtain that if $\bar{\theta}(y)$ is an optimal solution for the optimal design problem, then there exists a constant $C \in \mathbb{R}$ (coming from the Lagrange multiplier) such that

$$\bar{\theta}(x) = 1 \quad \text{if } \bar{U}(x)\bar{P}(x) - |\bar{U}(x)|^2 < C,$$

$$\bar{\theta}(x) \in (0, 1) \quad \text{if } \bar{U}(x)\bar{P}(x) - |\bar{U}(x)|^2 = C,$$

$$\bar{\theta}(x) = 0 \quad \text{if } \bar{U}(x)\bar{P}(x) - |\bar{U}(x)|^2 > C,$$

for a.e. $x \in [0, L]$, where U is the unique solution of

$$\bar{U}''(y) + \Omega^2 \sigma(\bar{\theta}(y)) \bar{U}(y) = 0, \quad y \in [0, L],$$

$$\bar{U}'(0) = \gamma, \quad \bar{U}'(L) = 0$$

and the adjoint state, Λ , is the unique solution of

$$[\Lambda']' + \Omega^2 \sigma(\bar{\theta}(y)) \Lambda = 2 (\Omega^2 \sigma(\bar{\theta}(y)) \bar{U} - \bar{U}''), \quad y \in [0, L],$$

$$\Lambda'(0) = -2\gamma, \quad \Lambda'(L) = 0.$$

The case $C = 0$ corresponds to the case without volume constraint studied on Theorem 3.1, so we assume now that $C \neq 0$. In this situation if we prove that the set

$$\{y \in [0, L] : \bar{U}(y)\Lambda(y) - |\bar{U}(y)|^2 = C\}$$

has zero measure, then we prove that problem (\tilde{P}) admits an optimal solution in the form of characteristic function, and consequently the original problem (P) admits at least one optimal solution. Unfortunately we have not been able to show this in general, however we claim that the condition

$$C < 0 \tag{6}$$

implies the result. We prove this by contradiction. Let us assume that the set

$$\bar{S} = \{y \in [0, L] : \bar{U}(y)\Lambda(y) - |\bar{U}(y)|^2 = C\}$$

has positive measure and that $C < 0$. For the sake of simplicity we assume that there exists an interval $[y_1, y_2] \subset \bar{S}$. First of all, notice that the function

$$\bar{V}(y) = \Lambda(y) - \bar{U}(y)$$

satisfies the equation

$$\bar{V}''(y) + \Omega^2 \sigma(\bar{\theta}(y)) \bar{V}(y) = 4\Omega^2 \sigma(\bar{\theta}(y)) \bar{U}(y), \tag{7}$$

a.e. $y \in [0, L]$. As \bar{U}, Λ are continuous functions,

$$\bar{U}(y)\Lambda(y) - |\bar{U}(y)|^2 = C$$

for all $y \in [y_1, y_2]$, and therefore

$$\bar{V}(y) = C/(\bar{U}(y)),$$

for all $y \in [y_1, y_2]$. Further $\bar{U} \in H^2(0, L)$, so that by the Sobolev embedding theorem, $\bar{U} \in C^1(0, L)$ and by the differentiability of Sobolev functions it is twice-differentiable almost every point, and its derivative is the usual one. Then computing \bar{V}'' and substituting in (7) we obtain that

$$2C \left((\bar{U}'(y))^2 + \Omega^2 \sigma(\bar{\theta}(y)) \bar{U}^2(y) \right) = 4\Omega^2 \sigma(\bar{\theta}(y)) \bar{U}^4(y),$$

what it is obviously a contradiction as the sign of the left hand side is negative and the one of the right hand side positive. This finishes the proof of the claim. The interest of this fact is just theoretical because, as far as the authors are able to say, we do not know a priori the sign of the constant C in order to establish a mathematical result. Although it shows there might be a lot situations in which the existence of at least one classical solution may be guaranteed. We give numerical evidence of this in the next section.

4 Numerical Approach and Examples

After the analysis carried out in the previous sections, this section is devoted to show several numerical simulations for our problem. The approach that we follow in the present work is to approximate directly the optimal solutions for the relaxed optimal design problem (\tilde{P}) by a gradient method. Of course this implies approximation of the states and adjoint states, i.e. solutions of the homogenized state equation (see Theorem 2.1) and the corresponding adjoint equation (see Theorem 2.2). We believe that this approach will certainly give good approximations of the optimal designs for problem (\tilde{P}) .

As we mentioned above the numerical algorithm used in this work is a gradient method.

More specifically, we have used a variable step-size descent method given by

$$\theta_{k+1} = \theta_k - \tilde{J}'(\theta_k)s_k,$$

where s_k is a small positive step such that $\tilde{J}(\theta_{k+1}) < \tilde{J}(\theta_k)$, $\tilde{J}'(\theta_k)$ is the gradient of the cost functional respect to the control, computed in Section 2, and the step s_k is looked for as the step that minimizes $\tilde{J}(\theta_{k+1}(s_k))$ (we use a descent method for computing s_k).

On the other hand, since the control must verify that $0 \leq \theta_{k+1} \leq 1$, then the updating formula can be easily compacted by

$$\theta_{k+1} = \max(0, \min(1, \theta_k - \tilde{J}'(\theta_k)s_k)),$$

being the same that we obtained by using a projected gradient method. Concerning the stopping criteria, it has been checked that the cost functional slightly changes its value and therefore, the L^∞ -norm of the difference $\tilde{J}(\theta_{k+1}) - \tilde{J}(\theta_k)$ reduces to zero, as we approximate to the minimum. All numerical examples belonging to this work have been obtained under the stopping criteria, $tol = 10^{-5}$, requiring for convergence about 15 iterations of the algorithm, on average. It has also been tested that very similar results are obtained when we run the same numerical examples with different mesh sizes.

For the sake of brevity, we illustrate our approach through two numerical examples.

In both of them, the design domain is an elastic rod of unit length subjected to time harmonic longitudinal excitation of amplitude γ in the left end. Material properties are $E_2 = 200$ GPa, $\rho_2 = 7.800$ kg/m³, and $E_1 = 70$ GPa, $\rho_1 = 2700$ kg/m³ corresponding

to steel and aluminum, respectively. In previous sections it has been possible to prove that optimized solutions are classical when the Young's modulus E keeps constant along the rod and to show that there are cases with E varying its value in which still there are classical solutions. In general, numerical simulations indicate that optimized layouts are classical, and therefore, there is no microstructure among the two phases we are mixing.

In the first example we have used the frequency value $\Omega = 50$ kHz. Even though we let the control θ vary continuously between 0 and 1 during the optimization process, we can see in Figure 1(a) that the optimized solution obtained is classical at the end of the process. Notice that this configuration make the wave damp itself in an optimized way, as it travels along the rod (see Figure 1(b)). Similar conclusions as before can be obtained in the second example for $\Omega = 200$ kHz (see Figure 1(c) and (d)). We can repeat the same numerical examples but keeping constant the Young's modulus, for instance, $E_2 = E_1 = 70$ GPa. The numerical results for $\Omega = 50$ kHz and 200 kHz (Figure 2) corroborate the existence of classical solutions such as it was analytically proved.

It is important to notice that (for this particular case $E_2 = E_1$) we initially tried with

a numerical approach based on optimality conditions. Although in some situations such algorithm gave us the same results as the ones obtained by the gradient method, the fact is that it did not enjoy good convergence properties in general, but when it did we got similar results.

In view of the numerical results we can state that, even though the design problem treated in this paper is rather different to the band-gap one, as we commented in the Introduction, the optimized finite structures (almost periodic in some cases) obtained (for minimum vibration energy) for a particular value of Ω , not avoid (like in the band-gap situation) but make rather difficult the wave can propagate itself along the rod only at that corresponding value of Ω .

To conclude, it is worth emphasizing that first, all this analysis makes sense for high frequency values (typically greater than 10 kHz), otherwise the optimal designs are “boring” because they would correspond to take only the material with higher ratio (E_i/ρ_i) , for $i = 1, 2$, and second, the same layouts we present here are obtained whenever the product Ωl keeps constant.

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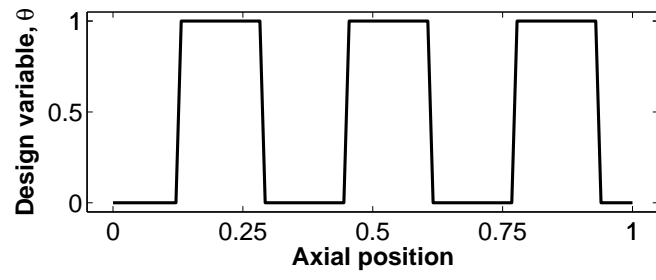
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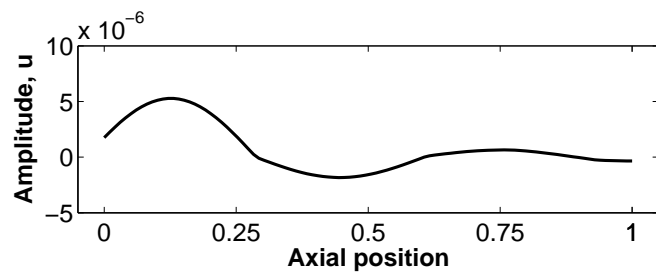
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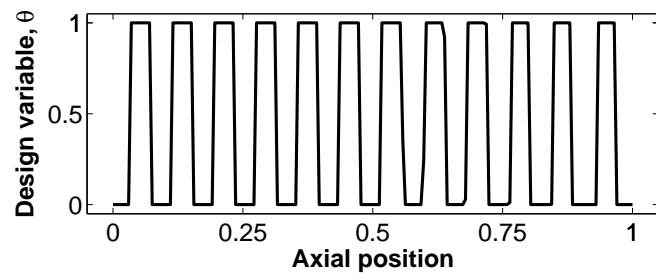
- Figure 1: Numerical results for $\Omega = 50$ kHz ((a) and (b)) and 200 kHz ((c) and (d)), with $E_2 > E_1$.
- Figure 2: Numerical results for $\Omega = 50$ kHz ((a) and (b)) and 200 kHz ((c) and (d)), with $E_2 = E_1$.



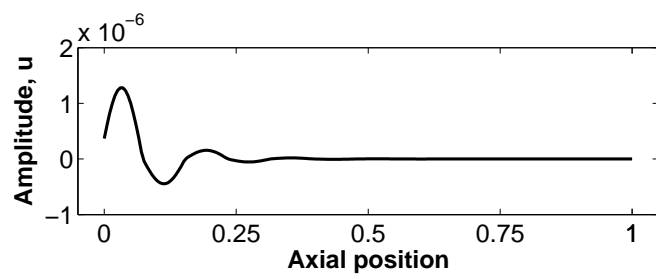
(a) Optimized design



(b) Wave amplitude

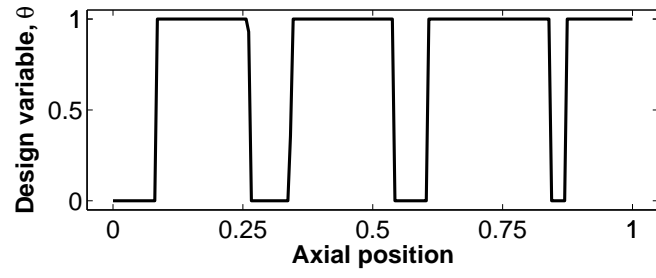


(c) Optimized design

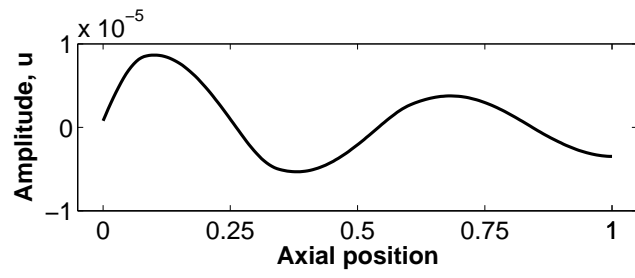


(d) Wave amplitude

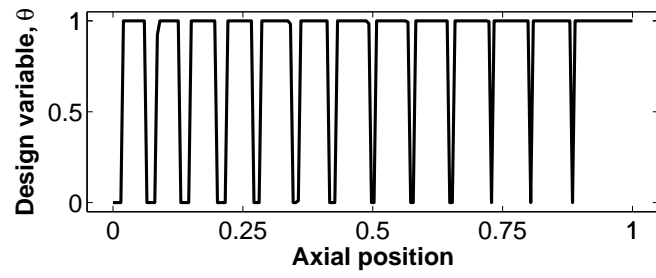
Figure 1:



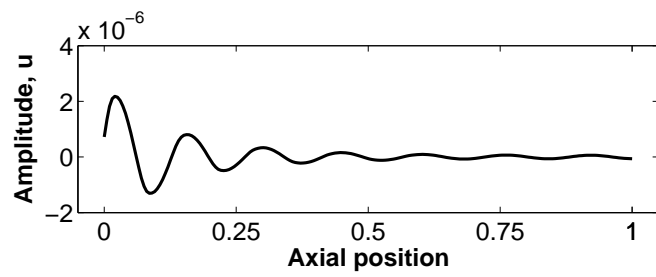
(a) Optimized design



(b) Wave amplitude



(c) Optimized design



(d) Wave amplitude

Figure 2: